

Problem set #2 Solutions

P. 209.

1) a)

$$G(\vec{r}; \vec{r}') = \frac{1}{\sqrt{\underbrace{(x-x')^2 + (y-y')^2}_{D^2} + (z-z')^2}}$$

$$G(x_i, y_i; x'_i, y'_i) = \int_{-Z}^Z G(\vec{r}, \vec{r}') d(z-z')$$

$$= \int_{-Z}^Z \frac{1}{\sqrt{D^2 + \xi^2}} d\xi$$

$$= \ln \left( \xi + \sqrt{D^2 + \xi^2} \right) \Big|_{-Z}^Z$$

$$= \ln \left( \frac{z + \sqrt{D^2 + z^2}}{-z + \sqrt{D^2 + z^2}} \right)$$

$$\stackrel{D \ll z}{=} \ln \left( \frac{z + z + D^2/z}{-z + z + D^2/z} \right) \sim \ln \left( \frac{2z^2}{D^2} \right)$$

 $D \ll z$ 

$$\boxed{G(\vec{r}; \vec{r}') = -\ln D^2 = -\ln[(x-x')^2 + (y-y')^2] = -\ln[\rho^2 + \rho'^2 - 2\rho\rho' \cos(\phi - \phi')]}$$

$$1) b) \quad \nabla^2 G^{2D}(\vec{r}, \vec{r}') = -4\pi \delta(\vec{r} - \vec{r}')$$

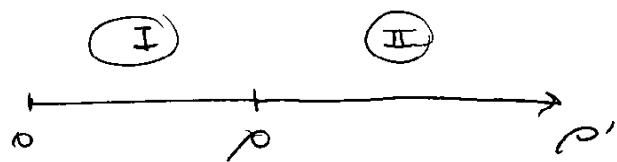
$$\left( \frac{1}{\rho'} \frac{\partial}{\partial \rho'}, \rho' \frac{\partial}{\partial \rho'} + \frac{1}{\rho'^2} \frac{\partial^2}{\partial \phi'^2} \right) \sum_{m=1}^{\infty} e^{im(\phi-\phi')} g_m(\rho; \rho') = -4\pi \delta(\rho - \rho') \frac{f(\phi)}{\rho'}$$

$$\begin{aligned} \sum_{m=1}^{\infty} \left( \frac{1}{\rho'} \frac{\partial}{\partial \rho'}, \rho' \frac{\partial}{\partial \rho'} + \frac{1}{\rho'^2} m^2 \right) g_m(\rho; \rho') e^{im(\phi-\phi')} &= -\frac{4\pi \delta(\rho - \rho')}{\rho'} \underbrace{\sum_{m=1}^{\infty} e^{im(\phi-\phi')}}_{2\pi} \\ &= \sum_{m=1}^{\infty} -\frac{4\pi \delta(\rho - \rho')}{\rho'} \frac{e^{im(\phi-\phi')}}{2\pi} \end{aligned}$$

$\Downarrow$  (uniqueness of  
Fourier series)

$$\boxed{\left[ \frac{1}{\rho'} \frac{\partial}{\partial \rho'}, \rho' \frac{\partial}{\partial \rho'} + \frac{1}{\rho'^2} m^2 \right] g_m(\rho; \rho') = -4\pi \frac{\delta(\rho - \rho')}{\rho'}}$$

1) c)



for regions I &amp; II

$$\left( \frac{1}{\rho}, \frac{\partial}{\partial \rho}, \rho, \frac{\partial}{\partial \rho}, -\frac{m^2}{\rho^{1/2}} \right) g_m^{I/II}(\rho) = 0$$

↓

$$\begin{cases} g_m^I = A\rho^{iml} + B\rho^{-lm} \\ g_m^{II} = C\rho^{ml} + D\rho^{-lm} \end{cases}$$

well behaved  
at 0 and  $\infty$  $g$  must be continuous at  $\rho' = \rho^*$ 

$$A\rho^{iml} = D\rho^{-lm}$$

$$\frac{A}{D} = \rho^{*-2ml}$$

and integrating the original equation

$$\frac{\partial}{\partial \rho'} g \Big|_{\rho=\rho^*}^{P+\epsilon} = -\frac{4\pi}{\rho}$$

$$D(-lm)\rho^{-lm-1} - A(m)\rho^{ml-1} = -4\pi/\rho$$

$$-D\rho^{-lm} = -\frac{4\pi}{m} + A(m)\rho^{ml}$$

$$-D = -\frac{4\pi}{(m)} \rho^{|m|} + A \rho^{2m}$$

$$-D = -\frac{4\pi}{(m)} \rho^{|m|} + D \cancel{\rho^{-2|m|}} \cancel{\rho^{2|m|}}$$

$$-2D = -\frac{4\pi}{(m)} \rho^{|m|}$$

$$D = \frac{2\pi \rho^{|m|}}{(m)} \quad ; \quad A = \frac{2\pi \rho^{-|m|}}{(m)}$$

final solution

$$\textcircled{I} \quad g_m(\rho' < \rho) = \frac{2\pi}{(m)} \rho^{-|m|} \rho'^{|m|} = \frac{2\pi}{(m)} \left( \frac{\rho'}{\rho} \right)^{|m|}$$

$$\textcircled{II} \quad g_m(\rho' > \rho) = \frac{2\pi}{(m)} \rho^{|m|} \rho'^{-|m|} = \frac{2\pi}{(m)} \left( \frac{\rho}{\rho'} \right)^{|m|}$$

$$g_m(\rho, \rho') = \frac{2\pi}{(m)} \left( \frac{\rho}{\rho'} \right)^{|m|}$$

↓

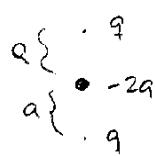
in the limit  $m=0$

$$g_0 = 2\pi \ln(\rho') + \text{constant}$$

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$$G(\rho; \rho'; \phi; \phi') = -\ln(\rho') + 2 \sum_{m=1}^{\infty} \frac{1}{m} \left( \frac{\rho}{\rho'} \right)^m \cos[m(\phi - \phi')]$$

2) a)



$$\Phi(x, y, z) = \frac{q}{\sqrt{x^2 + y^2 + (z-a)^2}} - \frac{2q}{\sqrt{x^2 + y^2 + z^2}} + \frac{q}{\sqrt{x^2 + y^2 + (z+a)^2}}$$

$$= \frac{q}{\sqrt{r^2 + a^2 - 2ar\cos\theta}} - \frac{2q}{r} + \frac{q}{\sqrt{r^2 + a^2 + 2ar\cos\theta}}$$

$$\begin{aligned} & \cancel{\frac{1}{\sqrt{r^2 + a^2 - 2ar\cos\theta}}} \cancel{\sqrt{r^2 - 2ar\cos\theta}} \\ & \cancel{\left[ r^2 - 2ar\cos\theta \right]^{\frac{1}{2}}} \\ & + \cancel{\frac{1}{8} \left[ r^2 - 2ar\cos\theta \right]^{\frac{3}{2}} (+8r\cos\theta)} \end{aligned}$$

Taylor series in  $a$ 

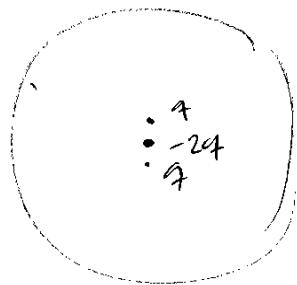
$$\left\{ \begin{aligned} & \left[ r^2 + a^2 \mp 2ar\cos\theta \right]^{\frac{1}{2}} \\ & \downarrow \\ & -\frac{1}{2} \left[ r^2 + a^2 \mp 2ar\cos\theta \right]^{\frac{3}{2}} (2a \mp 2r\cos\theta) \\ & \downarrow \\ & \left( -\frac{1}{2} \right) \left( -\frac{3}{2} \right) \left[ r^2 + a^2 \mp 2ar\cos\theta \right]^{-\frac{5}{2}} (2a \mp 2r\cos\theta)^2 + \left( -\frac{1}{2} \right) \left[ r^2 + a^2 \mp 2ar\cos\theta \right]^{\frac{3}{2}} (2) \end{aligned} \right.$$

$$\left[ r^2 + a^2 \mp 2ar\cos\theta \right]^{-\frac{1}{2}} \sim r^{-1} \pm r^{-3} 2r\cos\theta + \left[ \left( \frac{3}{4} \right) r^{-5} 4r^2 \cos^2\theta - \frac{3}{2} \right] \frac{a^2}{2}$$

$$\phi(x, y, z) = q \left( r^{-1} + r^{-3} 2r\cos\theta + r^{-3} (3\cos^2\theta - 1) \frac{a^2}{2} \right) - \cancel{2r^{-1}} + q \left( r^{-1} + r^{-3} 2r\cos\theta + r^{-3} (3\cos^2\theta - 1) \frac{a^2}{2} \right) =$$

$$\boxed{\Phi = \frac{2q}{r} r^{-3} \left( \frac{3\cos^2\theta - 1}{\alpha P_2(\cos\theta)} \right)} \leftarrow 4\text{-pole field.}$$

2(b) b) Turn-on sphere  
+  $q'$



+  $q'$

2 images + "image" of  $-2q$  at infinite = overall increase of potential.

$$\Phi_{\text{images}} = -q \left[ \frac{\frac{b/a}{(r^2 + \frac{b^2}{a^2} - 2rb^2 \cos\theta)^{1/2}} + \frac{b/a}{(r^2 + \frac{b^2}{a^2} - 2rb^2 \cos\theta)^{1/2}} - \frac{2}{b}}{b} \right]$$

using the same limit as in ↓

$$= +q \frac{1}{b} \frac{a^2 r^2}{b^4} (1 - 3\cos^2\theta)$$

this potential must be added to the previous one:

$$\boxed{\Phi_{\text{FINAL}} = 2Q \left( 1 - \frac{rs}{b^5} \right) P_2(\cos\theta)}$$

gaussian units.

~~BB/AA~~ 3a)

In the convention where all  $Q^{(l)}$  can be written as

$$Q_{x_i x_j \dots}^{(l)} = \underbrace{\int}_{\ell \text{ elements}} \underbrace{x_i x_j \dots}_{\ell} \rho(\bar{x}) d^3r \quad i, j, k \dots = 1, 2 \text{ or } 3$$

translating the origin:

$$\begin{aligned} \tilde{Q}_{x_i x_j \dots}^{(l)} &= \int x_i x_j \dots \rho(\bar{x} + \tilde{x}) d^3r \\ &= \int (x_i - \tilde{x}_i)(x_j - \tilde{x}_j) (x_k - \tilde{x}_k) \dots \rho(\bar{x}) d^3r \\ &\quad \text{expanding this product} \uparrow \\ &= \int x_i x_j x_k \dots \rho(\bar{x}) d^3r + \leftarrow \tilde{Q}_{x_i x_j \dots}^{(l)} \\ &+ -\tilde{x}_i \int (x_j x_k \dots) \rho(\bar{x}) d^3r \\ &+ -\tilde{x}_j \int (x_i x_k \dots) \rho(\bar{x}) d^3r \\ &\quad ; \\ &+ (\tilde{x}_i \tilde{x}_j \tilde{x}_k \dots) \int \rho(\bar{x}) d^3r \end{aligned}$$

combinations  
of  
 $\tilde{Q}^{(l')}$   
where  
 $l' < l$   
all equal to  $\emptyset$

so

$$= Q_{x_i x_j \dots}^{(l)}$$

$\text{Q.E.D.}$

~~3b)~~

$$\tilde{P}_i = \int_{\text{K}_i} \rho(\vec{r})$$

$$\begin{aligned}\tilde{P}_i &= \int (x_i - \tilde{x}_i) \rho(\vec{r}) = \int x_i \rho(\vec{r}) - \tilde{x}_i \int \rho(\vec{r}) \\ &= P_i - \tilde{x}_i Q\end{aligned}$$

$$\boxed{\tilde{P} = P - \tilde{R}Q}$$

$$r^2 = x_i x_i$$

$$(x_i - \tilde{x}_i)(x_i - \tilde{x}_i)$$

↓

$$x_i(x_i - \tilde{x}_i) - \tilde{x}_i(x_i - \tilde{x}_i)$$

$$x_i x_i - x_i \tilde{x}_i - \tilde{x}_i x_i + \tilde{x}_i \tilde{x}_i$$

$$x_i x_i - 2 x_i \tilde{x}_i + \tilde{x}_i \tilde{x}_i$$

$$r^2 - 2 \vec{r} \cdot \tilde{R} - \tilde{R}^2$$

$$Q_{ij} = \int (3 x_i x_j - r^2 \delta_{ij}) \rho(\vec{r}) d^3 r$$

$$\tilde{Q}_{ij} = \int [3(x_i - \tilde{x}_i)(x_j - \tilde{x}_j) - (r^2 - 2 \vec{r} \cdot \tilde{R} - \tilde{R}^2) \delta_{ij}] \rho(\vec{r}) d^3 r$$

$$\begin{aligned}&= \int (3 x_i x_j - 3 x_i \tilde{x}_j - 3 \tilde{x}_i x_j - 3 \tilde{x}_i \tilde{x}_j \\ &\quad - r^2 \delta_{ij} + 2 \vec{r} \cdot \tilde{R} \delta_{ij} + R^2 \delta_{ij}) \rho(\vec{r}) d^3 r\end{aligned}$$

$$\begin{aligned}&= \int (3 x_i x_j - r^2 \delta_{ij}) \rho(\vec{r}) - 3 \tilde{x}_i \int x_i \rho(\vec{r}) - 3 \tilde{x}_i \int x_j \rho(\vec{r}) \\ &\quad - 3 \tilde{x}_i \tilde{x}_j \int \rho(\vec{r}) d^3 r + 2 \vec{R} \cdot \tilde{R} \delta_{ij} \\ &\quad + 2 \tilde{R} \cdot \int \vec{r} \rho(\vec{r}) \delta_{ij} + R^2 \delta_{ij} \int \rho(\vec{r})\end{aligned}$$

$$\boxed{\tilde{Q}_{ij} = Q_{ij} - 3 \tilde{x}_j P_i - 3 \tilde{x}_i P_j - 3 \tilde{x}_i \tilde{x}_j Q + 2 \tilde{R} \cdot \tilde{P} \delta_{ij} + R^2 Q \delta_{ij}}$$

$$\tilde{Q} = Q - 3 \tilde{P} \cdot \tilde{R} - 3 \tilde{R} \otimes \tilde{P} - 3 R \otimes R Q + 2 \tilde{R} \cdot \tilde{P} \not{=} + R^2 Q \not{=} 1$$

3)c) consider  $\tilde{R} = \frac{P}{Q}$

for a general  $Q_{ij}$  there are 5 parameters  
an  $\tilde{R}$  only gives 3 parameters to play with  
so the system is over-constrained & in general  
has no sol.

$$4) a) W = \int \rho \Phi \quad \Phi = \Phi_0 + \partial_i \phi x_i + \frac{1}{2} \partial_i \partial_j \phi x_i x_j$$

$$W = \int \rho \Phi_0 + \int \rho \partial_i \phi x_i + \frac{1}{2} \int \rho \partial_i \partial_j \phi x_i x_j$$

$$= \Phi_0 \int \rho + \partial_i \phi \int \rho x_i + \frac{1}{2} \partial_i \partial_j \phi \int \rho x_i x_j + \dots$$

$$W(\vec{x}) = \phi(\vec{x}) q + \partial_i \phi(\vec{x}) p_i + \frac{1}{2} \partial_i \partial_j \phi \int \rho (\delta x_i x_j - \delta_{ij} r^2)$$

$$+ \cancel{\partial_{ij} \partial_i \partial_j \phi \int \rho r^2}$$

$$\cancel{\frac{\partial_i \partial_j \phi}{\nabla^2 \phi = 0}}$$

(\*)

$$W(\vec{x}) = \phi(\vec{x}) q + \partial_i \phi(\vec{x}) p_i + \frac{1}{6} \partial_i \partial_j \phi Q_{ij} + \dots$$

$$W(\vec{x}) = \phi(\vec{x}) - \vec{E} \cdot \vec{p} - \frac{1}{6} \partial_i E_j Q_{ij} + \dots$$

$$\vec{F} = -\vec{\nabla} W$$

$$\vec{F} = -\underbrace{\vec{\nabla} \phi(\vec{x})}_{\vec{E}_q} + \vec{\nabla}(\vec{E} \cdot \vec{p}) + \frac{1}{6} \vec{\nabla} \left[ \partial_i E_j Q_{ij} \right]$$

(\*) is still valid but now rotations can change the momenta (for example rotating the dipole moment)

$$q \rightarrow q \quad (\text{scalar})$$

$$\bar{p} \rightarrow \bar{p} + \hat{\mathbf{r}}_i \times \bar{p} \delta\theta$$

$$p_i \rightarrow p_i + \frac{\epsilon_{ijk} p_k \delta\theta}{\delta p_i}$$

$$Q_{ij} \rightarrow Q_{ij} + (\epsilon_{ikl} Q_{kj} + \epsilon_{jkl} Q_{ki}) \delta\theta$$

$$\overset{\text{SD}}{\delta W_1} = \left\{ \partial_i \phi \epsilon_{1lk} p_k + \frac{1}{6} \partial_i \partial_j \phi \left( \underbrace{\epsilon_{ikl} Q_{kj}}_{Q_{2j} \delta_{i3} - Q_{3j} \delta_{i2}} + \underbrace{\epsilon_{jlk} Q_{ki}}_{Q_{2i} \delta_{j3} - Q_{3i} \delta_{j2}} \right) \right\} \delta\theta$$

- N<sub>i</sub>

$$N_1 = - [\bar{p} \times \bar{\nabla} \phi]_1 - \frac{1}{6} \left\{ \partial_i \partial_j \phi Q_{2j} \delta_{i3} - \partial_i \partial_j \phi Q_{3j} \delta_{i2} + \partial_i \partial_j Q_{2i} \delta_{j3} - \partial_i \partial_j Q_{3i} \delta_{j2} \right\}$$

$$= - [\bar{p} \times \bar{\nabla} \phi]_1 - \frac{1}{6} \left\{ \partial_3 \partial_j \phi Q_{2j} - \partial_2 \partial_j \phi Q_{3j} + \partial_1 \partial_3 \phi Q_{2i} - \partial_1 \partial_2 \phi Q_{3i} \right\}$$

same.

$$N_1 = [\bar{p} \times \bar{E}] + \frac{1}{3} \left\{ \partial_3 E_j Q_{2j} - \partial_2 E_j Q_{3j} \right\}$$

This is an elegant way starting with  $\bar{N} = \int \rho \mathbf{r} \times \mathbf{E}(\mathbf{r})$  is also correct.

5) From 4)

$$5) W = -\frac{1}{6} (\partial_i E_j Q_{ij}) \rightarrow \text{by symmetry}$$

$$\bar{Q} = \begin{pmatrix} \epsilon_0 & 0 & 0 \\ 0 & \epsilon_0 & 0 \\ 0 & 0 & \epsilon_0 \end{pmatrix} \quad \text{"is like"} \quad \left. \begin{array}{l} \partial_x E_y = -\partial_y E_x \\ \text{and } \partial_x E_z = 0 \\ \partial_y E_x = 0 \\ \text{and } \partial_z E_x = 0 \\ \partial_z E_y = 0 \end{array} \right\} \begin{array}{l} \text{so off-diagonal } Q \\ \text{elements cancel each other.} \end{array}$$

$$W = -\frac{1}{6} \left[ \partial_x E_x \frac{\epsilon_0 Q}{2} - \partial_y E_y \frac{\epsilon_0 Q}{2} + \partial_z E_z \frac{\epsilon_0 Q}{2} \right] \quad \text{(Jackson is not very clear)}$$

$$= -\frac{e}{6} \left[ -(\partial_x E_x + \partial_y E_y) \frac{Q}{2} + \partial_z E_z Q \right]$$

$$\bar{E} \text{ is axially symmetrical} \Rightarrow \partial_x E_x = \partial_y E_y \quad \left. \begin{array}{l} \partial_x E_x = \partial_y E_y \\ \partial_x E_x + \partial_y E_y + \partial_z E_z = 0 \end{array} \right\}$$

$$\nabla \cdot \bar{E} = e \Rightarrow \partial_x E_x + \partial_y E_y + \partial_z E_z = 0 \quad \Downarrow$$

$$\partial_x E_x = \partial_y E_y = -\frac{\partial_z E_z}{2}$$

$$= -\frac{e}{6} \left[ -(-\partial_z E_z) \frac{Q}{2} + (\partial_z E_z) Q \right]$$

$$= -\frac{e}{6} \frac{3}{2} (\partial_z E_z) Q$$

$$W = -\frac{e}{4} Q (\partial_z E_z)$$

5) b)

$$Q = 2 \times 10^{-28} \text{ m}^2$$

$$\omega_h = 10 \text{ MHz}$$

$$W = -\frac{e}{q} \propto \left(\frac{\partial E}{\partial t}\right)_0$$

$$\left(\frac{\partial E}{\partial t}\right)_0 = -\frac{4W}{eQ} = \boxed{-0.08 \left(\frac{e}{4\pi\epsilon_0 a_0^3}\right)}.$$

5)c)

$$eQ = Q_{33} = \int (3z^2 - r^2) \rho(r) d^3x$$

$$= \rho \int_0^{2\pi} d\phi \int_{z=-b}^b dz \int_0^{\frac{a}{b} \sqrt{b^2-z^2}} (zz^2 - r^2) r dr$$

$$= \rho 2\pi \int_{-b}^b dz \left[ \frac{1}{2} z^2 r^2 - \frac{1}{4} r^4 \right] \Big|_0^{\frac{a}{b} \sqrt{b^2-z^2}}$$

$$= \rho 2\pi \int_{-b}^b dz \left\{ \frac{1}{2} z^2 \frac{a^2}{b^2} (b^2 - z^2) - \frac{1}{4} \frac{a^4}{b^4} (b^2 - z^2)^2 \right\}$$

$$= \rho 2\pi \int_{-b}^b \left[ \frac{z^2 a^2}{b^2} \right] - z^4 \frac{a^2}{b^2} - \frac{1}{4} \frac{a^4}{b^4} b^4 + \frac{a^4}{4b^4} 2b^2 z^2 - \frac{1}{4} \frac{a^4}{b^4} z^4$$

$$eQ = \rho 2\pi \frac{4}{15} a^2 b (b^2 - a^2) \leftarrow (\rho = Ze / \frac{4}{3} \pi a^2 b)$$

$$eQ = \frac{2}{3} Z(a+b)(a-b) \Rightarrow \boxed{\frac{a-b}{a} \approx 0.1} \leftarrow {}^{14}_6 \text{O} \leftarrow \text{nucleus} \leftarrow \text{fractional nuclear mass}$$